

Mathematical modelling of wind turbines



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The Betz Limit

The full wake

Turbulent far wake

$$\nu \neq 0$$

$$\nu \rightarrow 0$$

$$\nu = 0$$

How should we space wind turbines?

Stress analysis of blades

$$C_P = \frac{P_R}{P_W} = 59.259\%$$

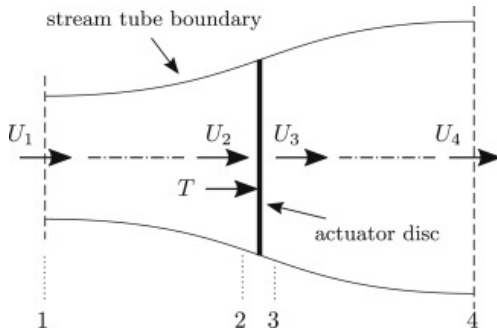


Figure: Simplified model used in the Betz limit.

We found using a similarity solution approach that,

$$w = F(\xi), \quad v_r = \frac{1}{\sqrt{z}}G(\xi), \quad \xi = \frac{r}{\sqrt{z}}.$$

Conserved quantities

$$D = z \int_0^\infty \xi F(\xi) d\xi, \quad J = z \int_0^\infty \xi F^2(\xi) d\xi.$$

$$v_z = \bar{v}_z + v'_z, \quad v_r = \bar{v}_r + v'_r.$$

$$v_z = U_0 - w$$

Kinematic Eddy viscosity

$$\nu_T = \ell^2(z) \left| \frac{\partial v_z}{\partial r} \right|. \quad (1)$$

Far wake approximation

1. We drop products of $\bar{w}$ and/or $\bar{v}_r$.
2. Retain the ν_T term.

$$\bar{w} = z^{-1}F(\xi), \quad \bar{v}_r = z^{-\frac{3}{2}}G(\xi), \quad \xi = \frac{r}{\sqrt{z}}.$$

The wake boundary and Prandtl's mixing length:

$$b(z) = \beta z^{\frac{1}{2}}, \quad \ell(z) = \ell_0 z^{\frac{3}{4}}.$$

$$D = U_0 \int_0^\beta \xi F(\xi) d\xi, \quad (2)$$

$$G(\xi) = -\frac{1}{2}\xi F(\xi), \quad (3)$$

$$F'(\xi) = \frac{\nu - \sqrt{\nu^2 + 2\ell_0^2 U_0 \xi F}}{2\ell_0^2}. \quad (4)$$

$$F(\xi) = \frac{U_0}{18\ell_0^2} \left(\beta^{\frac{3}{2}} - \xi^{\frac{3}{2}} \right)^2, \quad 0 \leq \xi \leq \beta,$$

where,

$$\beta = \left(\frac{140\ell_0^2 D}{U_0^2} \right)^{\frac{1}{5}}.$$

$$\bar{w} = z^{-\frac{2}{3}}F(\xi), \quad \bar{v}_r = z^{-\frac{4}{3}}G(\xi), \quad \xi = \frac{r}{z^{\frac{1}{3}}}.$$

The wake boundary and Prandtl's mixing length:

$$b(z) = \beta^* z^{\frac{1}{3}}, \quad \ell(z) = \ell_0^* z^{\frac{1}{3}}.$$

$$F(\xi) = \frac{U_0}{27\ell_0^{*2}} \left(\beta^{*\frac{3}{2}} - \xi^{\frac{3}{2}} \right)^2, \quad 0 \leq \xi \leq \beta^*,$$

where,

$$\beta^* = \left(\frac{210\ell_0^{*2}D}{U_0^2} \right)^{\frac{1}{5}}.$$

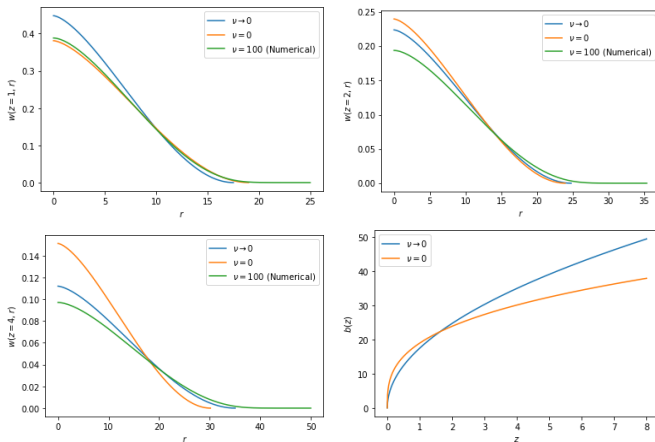
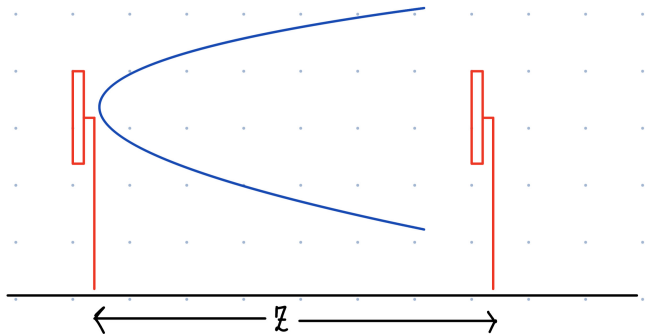


Figure: Numerical comparisons

How should we space wind turbines?

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$$\frac{\bar{w}_z^{max}}{U_0} \leq 1\%$$

$\nu \neq 0 :$

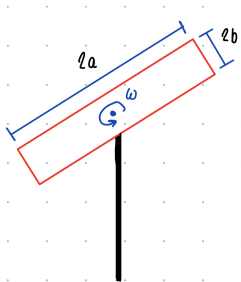
$$z \geq \frac{100F(0)}{U_0},$$

$\nu \rightarrow 0 :$

$$z \geq \frac{50\beta^3}{9\ell_0^2},$$

$\nu = 0 :$

$$z \geq \left(\frac{100\beta^{*3}}{27\ell_0^{*2}} \right)^{\frac{3}{2}}.$$



$$F_{net} = F + F^*,$$

where,

$$F^* = - \left(\frac{d\omega}{dt} \right)_R \times \underline{r} - \underline{\omega} \times (\underline{\omega} \times \underline{r}).$$

Boundary conditions:

1. Blade edge is traction free.
2. Tip of blades have no normal/shear force or moment.

Let $\bar{x} = \frac{x}{a}$ and $\bar{y} = \frac{y}{b}$.

Stresses:

$$\tau_{xx} = \frac{1}{2}\rho\omega^2 a^2 \left[1 - \bar{x}^2 + 2\sigma \left(\frac{b}{a} \right)^2 \left(\frac{1}{3} - \bar{y}^2 \right) \right],$$

$$\tau_{xy} = 0,$$

$$\tau_{yy} = \frac{1}{2}\rho\omega^2 b^2 (1 - \bar{y}^2).$$

Maximum normal stress

$$\tau^{max} = \frac{1}{2}\rho\omega^2 a^2 \left[1 + \frac{2\sigma}{3} \left(\frac{b}{a} \right)^2 \right]$$

This occurs at $(\bar{x}, \bar{y}) = (0, 0)$.

Maximum shearing stress

$$S_{max} = \frac{1}{4}\rho\omega^2 a^2 \left[1 - \frac{4\sigma}{3} \left(\frac{b}{a} \right)^2 \right]$$

This occurs at $(\bar{x}, \bar{y}) = (0, \pm 1)$.

The maximum displacement occurs at $(\bar{x}, \bar{y}) = (\bar{x}_0, 0)$,

$$\bar{x}_0 = \pm \sqrt{1 - \frac{\sigma}{3} \left(\frac{b}{a}\right)^2},$$

with magnitude,

$$u_x^{max} = \pm \frac{\rho \omega^2 a^3}{3E} \left[1 - \frac{\sigma}{3} \left(\frac{b}{a}\right)^2 \right]^{\frac{3}{2}}.$$

Note:

$$\left(\frac{b}{a}\right)^2 \approx 10^{-3}$$

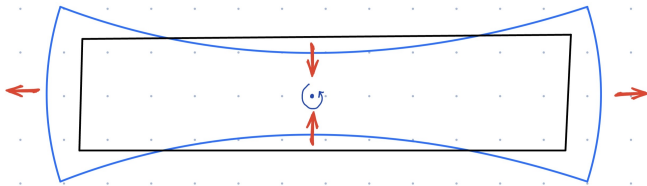


Figure: Deformed beam

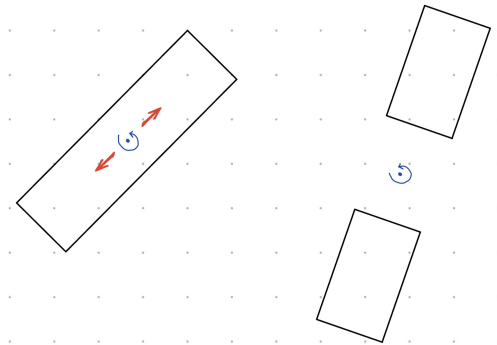


Figure: Normal stress fracture.

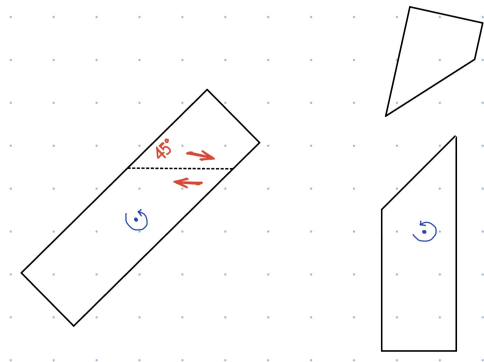


Figure: Shear stress fracture.

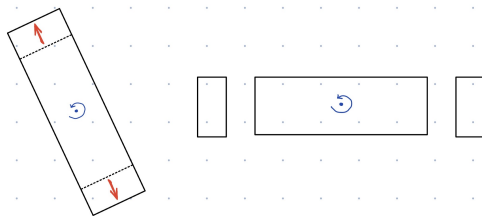


Figure: Extension stress fracture.